

## **Simulations of phonon-limited electron** mobility in II-VI semiconductors

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### INTRODUCTION

- Charge carrier mobility of pure crystalline materials is limited by the electron-phonon interaction.
- It is challenging to perform ab-initio prediction of mobility starting from the crystal structure of the material.
- Such studies were performed so far for only few representative semiconductors.
- Further studies are needed to further verify or improve the whole approach for mobility prediction.
- In this work we study II-VI semiconductors: ZnSe, CdTe, ZnTe and CdSe [1].

#### MOBILITY

The mobility is given as (for low electric field and small carrier concentration)

$$\mu_{xx} = e \frac{\sum_{mk} v_{x,mk}^2 \tau_{mk} \frac{\partial f_{mk}}{\partial \varepsilon_{mk}}}{\sum_{mk} f_{mk}}$$

- $f_{mk}$  is the Fermi Dirac factor
- $v_{x,mk}$  is the band velocity
- $\varepsilon_{mk}$  is the band energy
- $\tau_{mk}$  is the momentum relaxation time

 $\frac{1}{\tau_{mk}} = \frac{2\pi}{\hbar} \sum_{n\mu q} \left| \gamma_{nm}^{\mu}(\mathbf{k}, q) \right|^2 \begin{bmatrix} n_q^{\mu} \delta \left( \varepsilon_{mk} - \varepsilon_{nk+q} + \hbar \omega_q^{\mu} \right) + \\ (n_q^{\mu} + 1) \delta \left( \varepsilon_{mk} - \varepsilon_{nk-q} - \hbar \omega_q^{\mu} \right) \end{bmatrix} \left( 1 - \cos \theta_{\mathbf{k}, \mathbf{k} \pm q} \right)$ •  $\gamma_{nm}^{\mu}(\boldsymbol{k},\boldsymbol{q})$  are the e-ph coupling constants

### **MOBILITY - RESULTS**

Temperature dependence of mobility in ZnSe



#### **METHODS**

- Calculation on the coarse ( $\sim 8 \times 8 \times 8$ ) momentum grid:
- Density functional theory used to obtain the wave functions and energies of electronic states.
- Density functional perturbation theory used to obtain the electron-lattice coupling constants and the phonon modes.

Interpolation to a dense ( $\sim 100 \times 100 \times$ 100) momentum grid [2]:

- Electronic band structure
  - Construction of unitary matrices that connect the Bloch and the Wannier functions.
  - Representation of the electronic Hamiltonian in the Wannier basis.
  - **Diagonalization of this Hamiltonian** yields electronic energies at arbitrary electronic wave vector.
- Phonon dispersion / electron-phonon coupling constants

- $n_{a}^{\mu}$  are the Bose-Einstein factors
- $\hbar \omega_{\boldsymbol{q}}^{\mu}$  are the phonon energies

#### RESULTS

Electronic band structure of ZnSe obtained using PBE (black) and HSE06 (blue) functional



	ZnSe	ZnTe	CdSe	CdTe
$m_{ m eff} \exp$	0.137	0.117	0.119	0.090
$m_{ m eff}$ PBE	0.096	0.087	0.061	0.064
$m_{ m eff}$ HSE	0.14	0.13	0.13	0.13
$E_{\rm g}~({\rm eV})~{\rm exp}$	2.72	2.27	1.68	1.51
$E_{\rm g}~({\rm eV})~{\rm PBE}$	1.29	1.28	0.61	0.77
$E_{\rm g}~({\rm eV})~{\rm HSE}$	2.32	2.19	1.53	1.61

Phonon band structure of ZnSe 



points and dots – experiments from the literature full line - calculations

- PBE all calculations with PBE functional
- HSE band energies and high-frequency dielectric constants from the hybrid HSE functional

	ZnSe	ZnTe	CdSe	CdTe
$\varepsilon_{ m r}^{\infty}$ exp	5.9	6.9	6.2	7.1
$arepsilon_{ m r}^\infty$ PBE	7.26	8.86	7.72	8.62
$arepsilon_{ m r}^\infty$ HSE	6.15	7.45	6.13	7.01

### **COMPARISON OF** DIFFERENT MODELS



- Removal of the long-ranged part of dynamical matrices / electron-lattice coupling constants.
- Transformation of the short-ranged part to force-constant / Wannier representation.
- Evaluation of the short-ranged part on the dense momentum grid.
- Putting back the long-ranged part of dynamics matrices / electron-lattice coupling constants.
- Solving the eigenvalue problem of the dynamical matrices.
- Electron-phonon coupling constants in ZnSe



	Fröhlich	LR e-ph only	Full calc.
DFT			$\checkmark$
DFPT	X	$\checkmark$	$\checkmark$
Interpolation band structure	X		
Interpolation phonon modes	X		
Interpolation e-ph constants	X	X	
Accuracy	reasonable at room T	high	high

# REFERENCES

[1] N. Vukmirović, Phys. Rev. B 104, 085203 (2021). [2] F. Giustino et al, Phys. Rev. B 76, 165108 (2007).