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# Spectral functions of the Holstein polaron: exact and approximate solutions

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Center for the Study of Complex Systems  
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University of Belgrade  
Republic of Serbia*

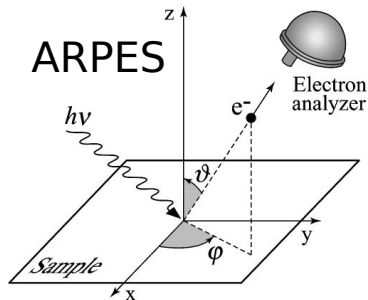
# Holstein model

Hamiltonian:

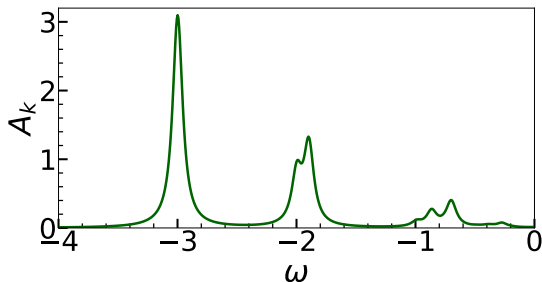
$$H = t_0 \sum_i c_i^y c_{i+1} + H:c: + g \sum_i n_i (a_i^y + a_i) + !_0 \sum_i a_i^y a_i$$

We calculate:

- q QP prop:  $m ; E_p$
- q **Spec. fun.**  $A_k(!)$
- q Im. time corr. fun.  $C_k( )$



Damascelli et al., Rev. Mod. Phys. **75**, 473 (2003)



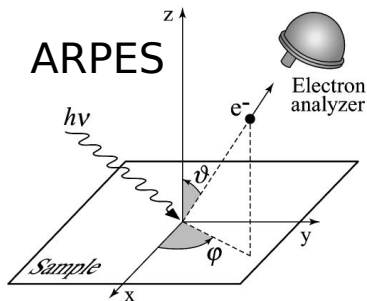
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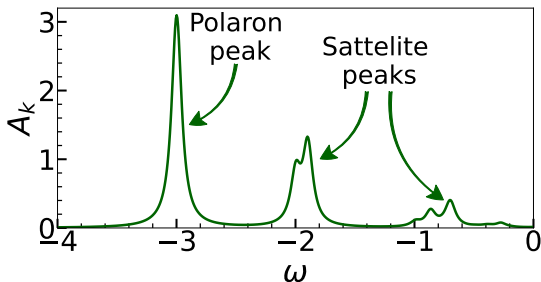
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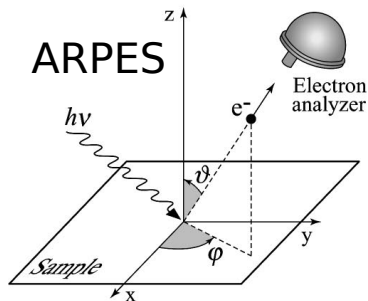
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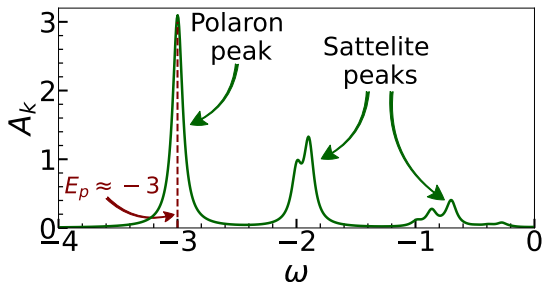
$$H = t_0 \sum_i c_i^y c_{i+1} + H:c: + g \sum_i n_i (a_i^y + a_i) + \sum_i \epsilon_0 a_i^y a_i$$

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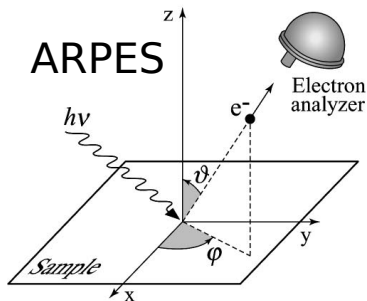
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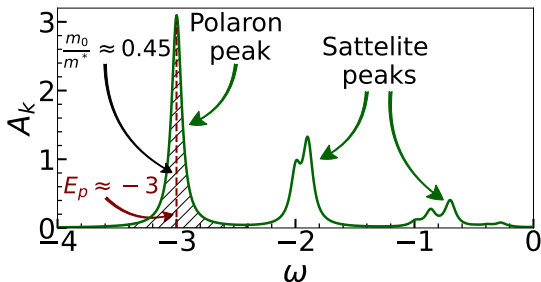
$$H = t_0 \sum_i c_i^y c_{i+1} + H:c: + g \sum_i n_i (a_i^y + a_i) + \frac{1}{2} \sum_i a_i^y a_i$$



Damascelli et al., Rev. Mod. Phys. **75**, 473 (2003)

We calculate:

- q QP prop:  $m^*$ ;  $E_p$
- q **Spec. fun.**  $A_k(\omega)$
- q Im. time corr. fun.  $C_k(\omega)$



# Holstein model

Hamiltonian:

$$H = t_0 \sum_i c_i^y c_{i+1} + H:c: + g \sum_i n_i a_i^y + a_i + \sum_i \epsilon_0 a_i^y a_i$$

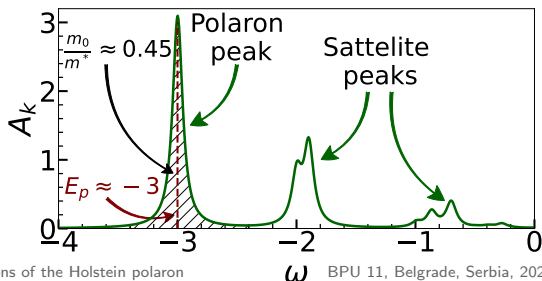
$$C_k(\omega) = \sum_{l=0}^{\infty} A_k(l) e^{i\omega \tau} d^l$$

$$A_k(l) = \frac{1}{l!} \text{Im} G_k(l)$$

$$G_k(l) = \frac{1}{i^l + 2t_0 \cos k} \frac{1}{k(i)}$$

We calculate:

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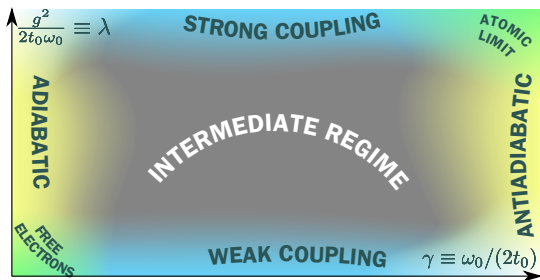
$$H = t_0 \sum_i c_i^y c_{i+1} + H:c:$$

$$g \sum_i n_i a_i^y + a_i$$

$$+ \sum_i \frac{1}{2} \omega_0 a_i^y a_i$$

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Mitrić et al., Phys. Rev. Lett. **129**, 096401 (2022)

## q Dynamical mean-field theory (DMFT)

- q Self-consistent Migdal approximation (SCMA)
- q Hierarchical equations of motion (HEOM)
- q Quantum Monte Carlo (QMC)
- q Exact diagonalization (ED)

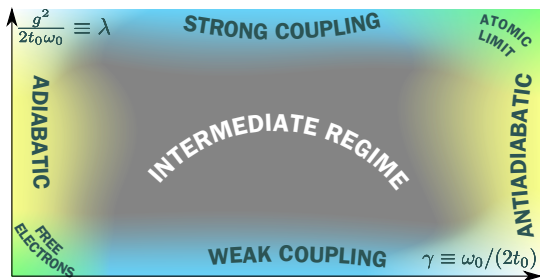
# Holstein model

Hamiltonian:

$$H = \sum_i t_0 c_i^\dagger c_{i+1} + H:c: \\ + \sum_i g \times^i n_i a_i^\dagger + a_i \\ + \sum_i !_0 \times^i a_i^\dagger a_i$$

We calculate:

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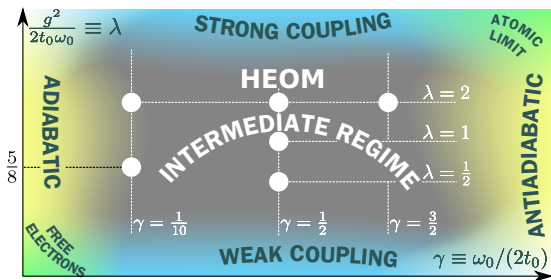
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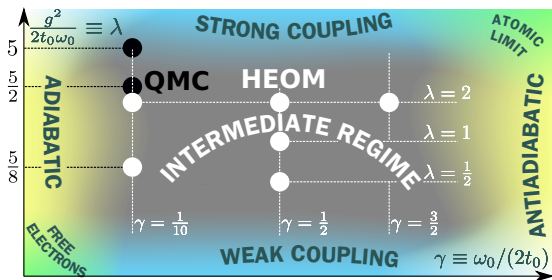
# Holstein model

Hamiltonian:

$$H = \sum_i t_0 c_i^\dagger c_{i+1} + H:c: + \sum_i g X_i^j n_i a_i^\dagger + a_i + \sum_i !_0 X_i^i a_i^\dagger a_i$$

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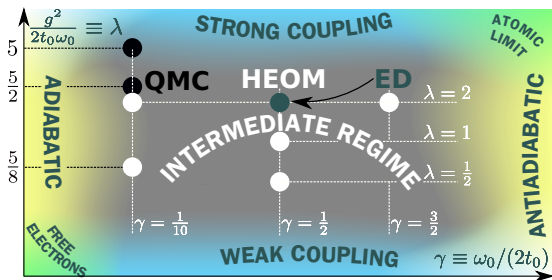
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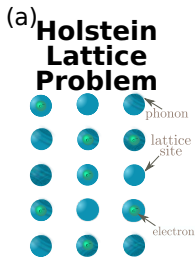


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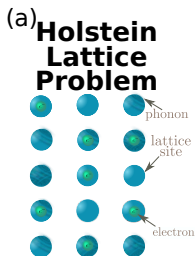
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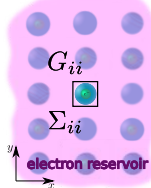
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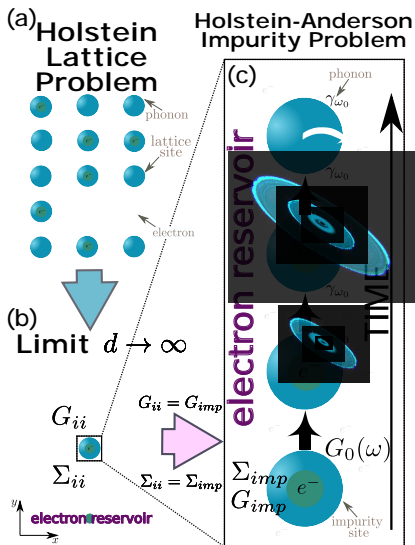
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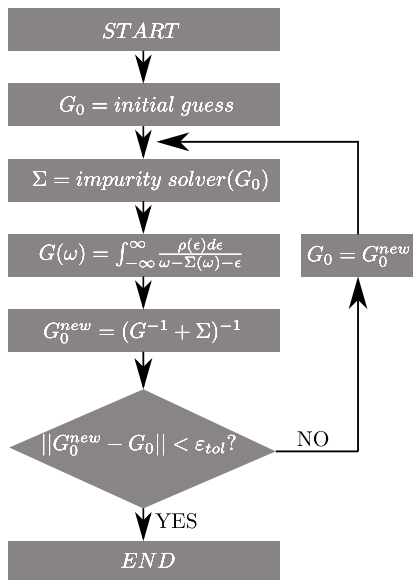
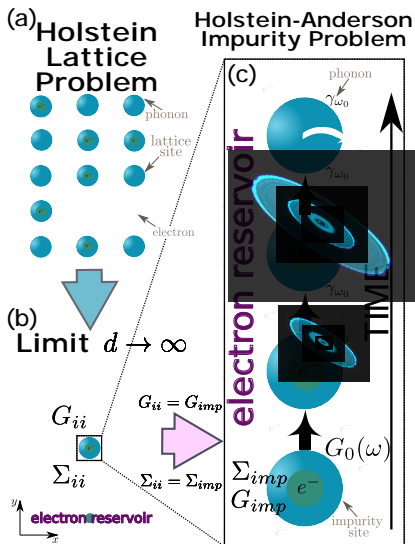
(b) **Limit**  $d \rightarrow \infty$



# Dynamical mean-field theory(DMFT)



# Dynamical mean-field theory(DMFT)



# Results: Spectral functions

Parameter regime:

Left Fig:

q (a){(d) Intermediate

q (e){(f) Strong

q (g){(h) Atomic

q Weak

# Results: Spectral functions

Parameter regime:

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# Results: Spectral functions

Parameter regime:

Left Fig:

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Bottom Fig:

q Weak

# Results: 1D Quasiparticle properties

## q Effective mass: DMFT vs DMRG

E. Jeckelmann and S. R. White, Phys. Rev. B 57, 6376 (1998)

## q Ground state energy: DMFT vs Global-local variational approach

A. H. Romero, D. W. Brown, and K. Lindenberg, J. Chem. Phys. 109, 6540 (1998)

Mitrić et al., Phys. Rev. Lett. 129, 096401 (2022)

# Higher dim. QP prop. and im. time corr. func.

- q Effective mass: DMFT vs Continuous-time QMC

Kornilovitch, Phys. Rev. Lett. 81, 5382 (1998)

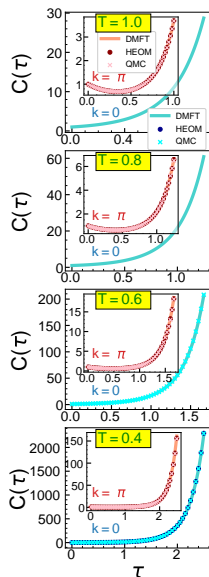
- q Imaginary time correlation functions ( $t_0 = 1, g = \sqrt{2}$ )

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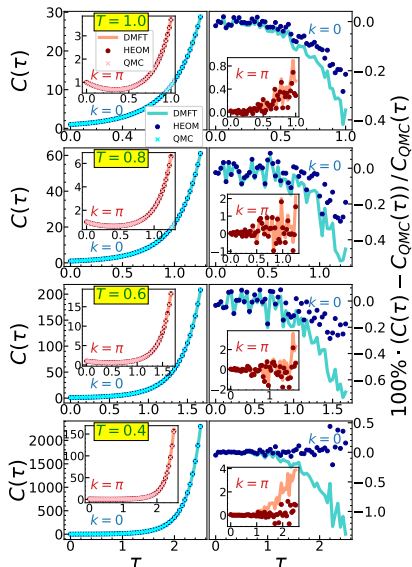
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q Imaginary time correlation functions ( $t_0 = 1, g = \frac{1}{2}$ )



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# Epilogue

## Take-Home Message

DMFT gives highly accurate and numerically cheap approximate solutions of the Holstein model (nonlocal correlations in Holstein model are small)

- q Covers the whole parameter regime which consists of  $t_0; g; T; k$  and number of dimensions.
- q Applicable both at the thermodynamic limit and on a chain with finite length.
- q Drawback: Gives local  $\chi_{\mathbf{k}}$ .
- q Possible applications to other systems:
  - q Extended DMFT for the non-local electron-phonon interaction.
  - q Statistical DMFT for the systems with disorder.
- q all details in P. Mitric, V.Jankovic, N. Vukmirovic and D. Tanaskovic, Phys. Rev. Lett. **129**, 096401 (2022)